# Gravitational radiation from shock wave collisions in higher dimensions 

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## IV Black Hole Workshop

In collaboration with Flávio Coelho, Carlos Herdeiro \& Carmen Rebelo

## Transplanckian collisions @ speed of light: Motivation

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## Evidence for classical BH in transplanckian scattering

- Numerical relativity in 4 and higher dimensions
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## LHC pp collisions well above 1 TeV !



The BH is assumed to decay through Hawking evaporation.
(?) As we approach $M_{4+n} \Rightarrow$ unknown quantum gravity effects!
First bound's rely on "bad" knowledge of gravitational radiation
CMS collaboration arXiv:1012.3357
ATLAS-CONF-2011-065

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(1) Aichelburg-Sexl shock waves

- Definition \& Physical interpretation
- Superposition and apparent horizons
(2) Shock wave collisions in D-dimensions
- Perturbative set up to determine future development
- The first order calculation \& radiation extraction
(3) Conclusions and Outlook


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## The Aichelburg-Sexl ultraboost

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d s^{2}=-d u d v+d \rho^{2}+\rho^{2} d \Omega_{D-3}^{2}+\kappa \Phi(\rho) \delta(u) d u^{2}
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(u, v)=(t-z, t+z)
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Flat region I

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## Basic properties of a single shock wave

- Solution of Einstein's equations, point source $P^{\mu}=E n^{\mu}$

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T^{\mu \nu}=E \delta(u) \delta^{(D-2)}\left(x^{i}\right) n^{\mu} n^{\nu} \quad, \quad n^{\mu} n_{\mu}=0
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- On the shock we have a profile

- Riemann tensor singular on the shock
- Null geodesics and tangent vectors are ciscontinuous
- No difference if we smear \& quantum corrections small
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## Superposition of two shock waves



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## Constraints due to the apparent horizon

- Apparent horizon area $\Rightarrow$ lower bound on $M_{\text {trapped }}$


Frost, Gaunt, MOPS, Casals, Dolan, Parker, Webber arXiv:0904.0979
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- $\Rightarrow$ Upper bound on the amount of gravitational radiation
- News function result 16.3\% consistent with $14 \pm 3 \%$ NGR


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## Strong shock vs weak shock

D'Eath and Payne found a trick $\Rightarrow$ perturbative approach
(ㄱ) Perform a large boost along $+z$ with velocity $\beta \equiv \tanh \alpha$
(2) Energy parameter of right/left moving shock $\kappa \rightarrow e^{ \pm \alpha} \kappa$
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$$
g_{\mu \nu}=\nu^{\frac{2}{D-3}}\left[\eta_{\mu \nu}+\frac{\lambda}{\nu} h_{\mu \nu}^{(1)}+\left(\frac{\lambda}{\nu}\right)^{2} h_{\mu \nu}^{(2)}\right]
$$

## Validity \& Physical interpretation

Note: $h_{\mu \nu}^{(i)} \sim\left[\frac{1}{\rho^{0-2}}(\sqrt{2} v-\Phi) \theta(\sqrt{2} v-\Phi)\right]^{i}$



## Future light cone of the collision

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u=0, \quad v=\frac{\Phi(\rho)}{\sqrt{2}}
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## Perturbative expansion

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- Fix gauge order by order (de Donder)
- Obtain decoupled sets of wave equations with source


Note: The source is zero for $i=1$ !

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## Causal structure of the background \& formal solution



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## Integration limits and ray analysis 1



Note: Point at the axis is a blind spot $\leftarrow$ destructive interference

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## Ray analysis off axis



## The first order result - radiation extraction

- Radiative components $h_{i j}=\delta_{i j} D(u, v, \rho)+\Delta_{i j}(x) E(u, v, \rho)$
- At first order ( $i=1$ ) only $E(u, v, \rho)$
- In our paper $\rightarrow$ Landau-Lifshitz pseudo-tensor

Yoshino and Shibata, arXiv:0907.2760

- Recently shown equivalent to news function in higher D Tanabe, Kinoshita and Shiromizu, arXiv:1104.0303



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\begin{aligned}
E_{\text {radiated }} & =\int d t \int_{S^{D-2}} \frac{d E n e r g y}{d S d t} \\
& \simeq \frac{\Omega_{D-3}}{32 \pi G_{D}} \lim _{\hat{\theta} \rightarrow 0, r \rightarrow \infty}\left(r^{2} \rho^{D-4} \int h^{i j}{ }_{, v} h_{i j, v} d t\right) \\
\frac{E_{\text {radiated }}}{E_{C M}} & \rightarrow \frac{1}{8} \frac{D-2}{D-3} \lim _{\hat{\theta} \rightarrow 0, r \rightarrow \infty}\left(\int\left(r \rho^{\frac{D-4}{2}} E_{, v}\right)^{2} d t\right)
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## Reduced Wave forms




## Extracting the axis limit




| D | 4 | 5 | 6 | 7 | 8 | 9 | 10 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| AH bound (\%) | 29.3 | 33.5 | 36.1 | 37.9 | 39.3 | 40.4 | 41.2 |
| First order(\%) | 25.0 | - | 33.3 | - | 37.5 | - | 40.0 |

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(1) Transplanckian collision in $D>4$ are important for Black Holes at the LHC.
(2) We have generalised the D'Eath \& Payne calculation and found a first order estimate that indicates a smaller emission of gravitational radiation than AH bounds.
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The main challenges are:
i) computational (numerically more intensive)
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4 Second order calculation (which gives good results in 4D) on the way. The main challenges are: i) computational (numerically more intensive) ii) extracting the first angular correction from numerics

## Conclusions

(1) Transplanckian collision in $D>4$ are important for Black Holes at the LHC.
(2) We have generalised the D'Eath \& Payne calculation and found a first order estimate that indicates a smaller emission of gravitational radiation than AH bounds.
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Open questions in progress

- How to solve odd $D$ ? (non-itegrable tails)
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Thanks for your attention! Questions?

BACKUP

The Standard Model - Particle content

## The Standard Model - Particle content

"Low" energy degrees of freedom (after symmetry breaking):

- 1 Higgs particle $(s=0)$,
- 3 families of leptons and 3 of quarks
- 1 non-abelian SU(3) c gluon field, 3 massive vector bosons, 1 neutral $U(1)$ Maxwell field


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The Standard Model - Interactions


The hierarchy problem: SM vs Gravity
The action for gravity coupled to matter is

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| Operator type | Couplings | at $E \sim 1 \mathrm{TeV}$ |
| :---: | :---: | :---: |
| $T_{\alpha \beta} h^{\alpha \beta}$ | $E / M_{4}$ | $10^{-16}$ |
| SM Interactions | $\sim e, g_{Q C D}, \frac{m_{H}}{v}, \frac{v}{E}, \frac{m_{f}}{v}$ | $O\left(10^{-6}\right)-O(1)$ |

## Solving the hierarchy problem with Extra Dimensions

$$
\mathbf{M}_{\mathbf{4}} \sim 10^{16} \mathbf{M}_{\mathrm{EW}}
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Hierarchy due to taking the scale for new physics from gravity (mesoscopic) rather than the electroweak scale (microscopic). The ADD solution: Assume $\mathbf{M}_{\mathrm{EW}}$ is more fundamental.
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- Predicts deviations from Newtonian gravity as we approach short distances.
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$\Rightarrow$ Translates as a bound on $M_{4+n}$.

## Bounds on extra dimensions

| $\mathbf{M}_{\mathbf{4}}^{\mathbf{2}}=\mathrm{R}^{\mathrm{n}} \mathbf{M}_{(4+\mathbf{n})}^{2+\mathbf{n}}$ | R in $\mu \mathrm{m}(\mathbf{n}=2)$ | $\mathbf{M}_{\mathbf{4 + \mathbf { n }}} \sim 1 \mathrm{TeV} \mathrm{OK}$ |
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- SM on a 4D brane of thickness $L \lesssim(1 \mathrm{TeV})^{-1} \sim 10^{-13} \mu \mathrm{~m}$

To avoid bounds from Electroweak precision and fast proton decay. Quarks and leptons may have to be on sub-branes for $L \lesssim(1 \mathrm{TeV})^{-1}$.

- All SM particles propagating on a single brane. Good approximation if process occurs at large scales compared to $L$.


## Gravity becomes strong above $M_{D} \sim 1 \mathrm{TeV}$

At short distances gravity is higher dimensional

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\Rightarrow \sqrt{\alpha_{G}} \sim \frac{E}{M_{4}} \rightarrow \frac{E}{M_{4+n}} \sim \frac{E}{1 \mathrm{TeV}}
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S. B. Giddings and S. D. Thomas, hep-ph/0106219
S. Dimopoulos and G. Landsberg, hep-ph/0106295

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Hoop conjecture $\Rightarrow \sigma_{\text {disk }} \sim \pi \mathbf{r}_{\mathbf{S}}^{2}, \quad \mathbf{r}_{\mathbf{s}}=\frac{C_{n}}{M_{4+n}}\left(\frac{\sqrt{s}}{M_{4+n}}\right)^{\frac{1}{n+1}}$
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## Evidence for classical BH in transplanckian scattering

- Numerical relativity in 4 and higher dimensions
U. Sperhake, V. Cardoso, F. Pretorius, E. Berti, J. Gonzalez, arXiv:0806.1738 b=0
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The BH is assumed to decay through Hawking evaporation.
(?) As we approach $M_{4+n} \Rightarrow$ unknown quantum gravity effects!
First bounds rely on bad knowledge of gravitational radiation
CMS collaboration arXiv:1012.3357
ATLAS-CONF-2011-065

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If $\Lambda_{\text {cutoff }} \sim M_{4} \sim 10^{16} \mathrm{TeV} \Rightarrow$ fine tuning of $\sim 10^{-16}$

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Also quantum gravity approximations indicate small corrections:
T. Banks and W. Fischler, hep-th/9906038
S. N. Solodukhin, hep-ph/0201248
S. D. H. Hsu, hep-ph/0203154

## Transient period

- During formation we should have an asymmetric BH with electric and gravitational multipole moments.
$\rightarrow$ Distorted geometry.
- The time for loss of multipoles is $\mathrm{r}_{s}$ (natural units).
- We will look noxt into the Hawking decay and realise that the typical timescale there is

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$$

We assume a quick loss of asymmetries
$\Rightarrow \mathrm{BH}$ settles down to a stationary axisymmetric solution.

## Apparent horizon



