



VII Black Holes Workshop

# Light-rings as observational evidence for event horizons

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# Introduction & Motivation

- Many massive stars are unstable against gravitational collapse.
- Exotic stars: boson stars and gravastars.
- Gravitational wave physics.
- Very compact spinning objects can be unstable: **Ergoregion instability**.
- Ultracompact objects: No horizons and **Light-rings** (Comins and Schutz, Proc. R. Soc. Lond. A **364**, 211).
- Can **ultracompact objects** be unstable? (Keir, 1404.7036)
- Here we study the linear modes of **ultracompact objects**.

# Ultracompact objects

Spherical **ultracompact** objects are described by

$$ds^2 = -f(r)dt^2 + B(r)dr^2 + r^2d\Omega_2^2. \quad (1)$$

The radial equation for null geodesics in the above spacetime is given by

$$Bf\dot{r}^2 = E^2 - V_{\text{geo}} \equiv E^2 - L^2 \frac{f}{r^2}. \quad (2)$$

The maxima and minima points of  $V_{\text{geo}}$  corresponds to unstable and **stable** geodesics.

# Eikonal limit of the modes

- Schwarzschild spacetime: Unstable LR. (Cardoso *et al.*, 0812.1806)  
 $\Re(\omega) \rightarrow$  Angular frequency of the circular null geodesic.  
 $\Im(\omega) \rightarrow$  Instability timescale of the circular null geodesic (Lyapunov exponents).
  
- Ultracompact objects: Stable LR.  
 $\Re(\omega) \rightarrow$  Angular frequency of the *stable* circular null geodesic.  
 $\Im(\omega) \rightarrow$  Becomes arbitrarily small.

# Ultracompact models: Constant density stars and gravastars

- Constant density stars: Idealized (toy) model

$$f(r) = \frac{1}{4R^3} \left( \sqrt{R^3 - 2Mr^2} - 3R\sqrt{R - 2M} \right)^2, \quad (3)$$

$$B(r) = \left( 1 - \frac{2Mr^2}{R^3} \right)^{-1}. \quad (4)$$

- Gravastars: Devised to mimic black holes (Mazur and Mottola, gr-qc/0109035)

$$f(r) = B(r)^{-1} = 1 - \frac{2M}{R} \frac{r^2}{R^2}. \quad (5)$$

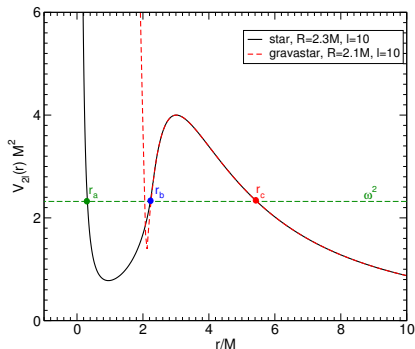
# Perturbations

Various classes of perturbations are described by a master equation

$$\left[ \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial r_*^2} + V_{sl}(r) \right] \Psi(r, t) = 0, \quad (6)$$

where  $\frac{\partial^2}{\partial r_*^2} = \frac{f}{B} \frac{\partial^2}{\partial r^2} + \frac{f}{2B} \left( \frac{f'}{f} - \frac{B'}{B} \right) \frac{\partial}{\partial r}$  and

$$V_{sl}(r) = f \left[ \frac{l(l+1)}{r^2} + \frac{1-s^2}{2rB} \left( \frac{f'}{f} - \frac{B'}{B} \right) + 8\pi(p_{\text{rad}} - \rho)\delta_{s2} \right]. \quad (7)$$



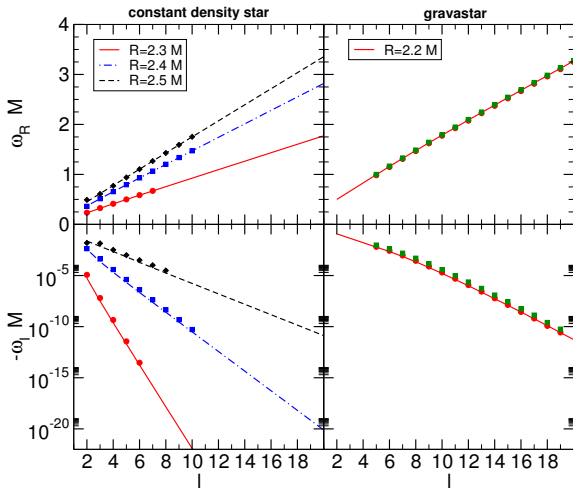
$$\int_{r_a}^{r_b} \frac{dr}{\sqrt{f/B}} \sqrt{\omega_R^2 - V_{sl}(r)} = \pi (n + 1/2), \quad (8)$$

$$\omega_I = -\frac{1}{8\omega_R\gamma} e^{-\Gamma}, \quad (9)$$

$$\omega \sim \Omega_{\text{LR}2} l - i b e^{-cl} \quad l \gg 1. \quad (10)$$



# Numerical results: Spectrum of linear perturbations

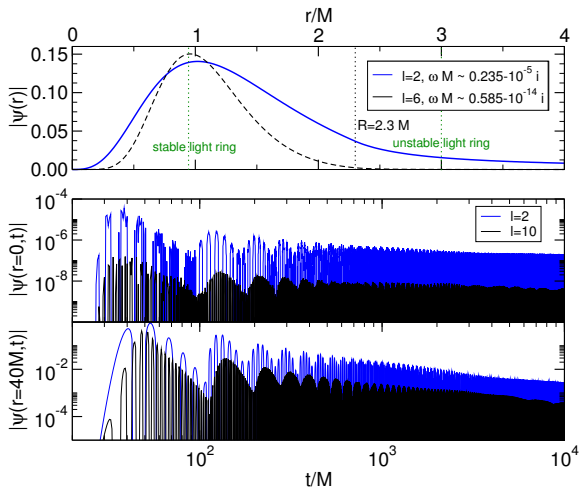


# Numerical results: Time evolution of wavepackets

We evolve an initial Gaussian wavepack in the star spacetime:

$$\dot{\Psi}(0, r) = \exp \left[ -\frac{(r + 2 \log (r - R) - r_0)^2}{\sigma^2} \right]. \quad (11)$$

# Numerical results: Time evolution of wavepackets



# Ergoregion instability (Comins and Schutz, Proc. R. Soc. Lond. A **364**, 211)

If we set rotation on, we have

$$ds^2 = -F(r)dt^2 + B(r)dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta (d\phi - \varpi(r)dt)^2, \quad (12)$$

and the scalar field equation is

$$\psi'' + m^2 \frac{B}{F} (\bar{\omega} + V_+) (\bar{\omega} + V_-) \psi = 0, \quad (13)$$

where

$$V_{\pm} = -\varpi \pm \frac{\sqrt{F}}{r}. \quad (14)$$

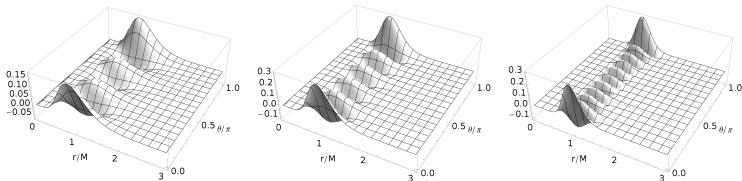
# Possible nonlinear outcomes

There are some possible outcomes from the nonlinear regime:

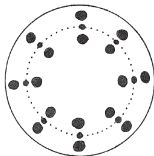
- **Other dissipative mechanisms become relevant:** No collapse.
- **Nonlinear effects become relevant:** Collapse or fragmentation.

# Possible nonlinear outcomes

Formation of small BHs (Bizon and Rostworowski, 1104.3702v5), (Okawa *et al.*, 1409.0533v2)



“Boiling fluid” (Lehner and Pretorius, 1006.5960)



- We show evidences that **any object with a light ring are BHs**.
- This is based on the fact that these objects possesses **long-living modes** at linear level.
- **Two outcomes** are possible: Decreasing of compactness or BH formation.
- For BH formation: **Nonlinear instabilities and weakly turbulence**.
- Turning on rotation: **Ergoregion instability (linear level)**.

# Acknowledgments



<http://centra.tecnico.ulisboa.pt/network/grit/>



<http://ppgf.ufpa.br/>