

Growth and decay of scalar clouds outside (rotating) black holes



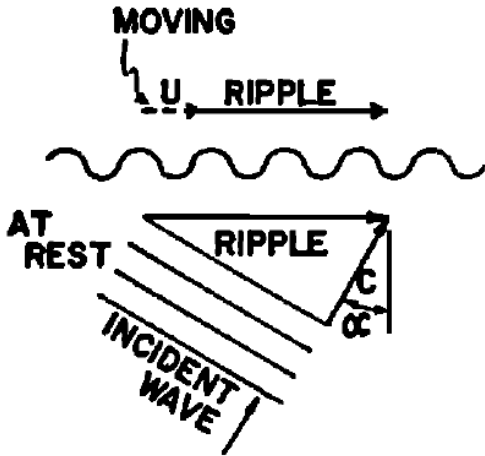
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Berti, Brito, Carucci, Gualtieri, Ishibashi, Okawa, Pani, Sotiriou, Sperhake, Witek

* * *

Brito, Cardoso, Pani, Superradiance in Black Hole Physics 2015

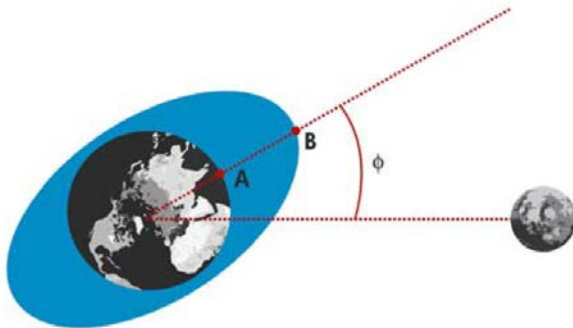
Friction & superradiance



Ribner, J. Acous. Soc. Amer. 29 (1957)

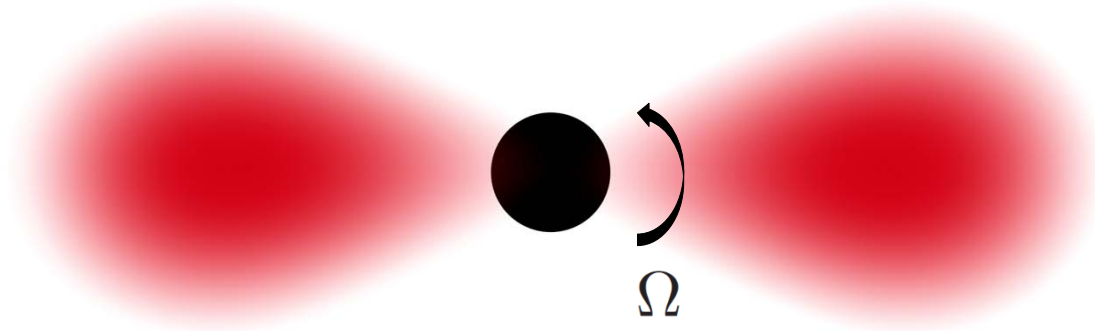


Tamm & Frank, Doklady AN SSSR 14 (1937)



G. H. Darwin, Philos. Trans. R. Soc. London 171 (1880)

$$\Phi \sim e^{-i\omega t + im\phi} \rightarrow (\text{Angular}) \text{ phase velocity} = \frac{\omega}{m}$$



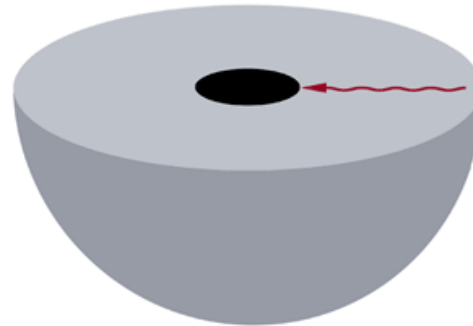
$$\omega < m\Omega$$

Zel'dovich, Pis'ma Zh. Eksp. Teor. Fiz. 14 (1971)

Black holes and superradiance

Friction built-in through one-way membrane (horizon)

Can construct unstable states by forcing wave to bounce back

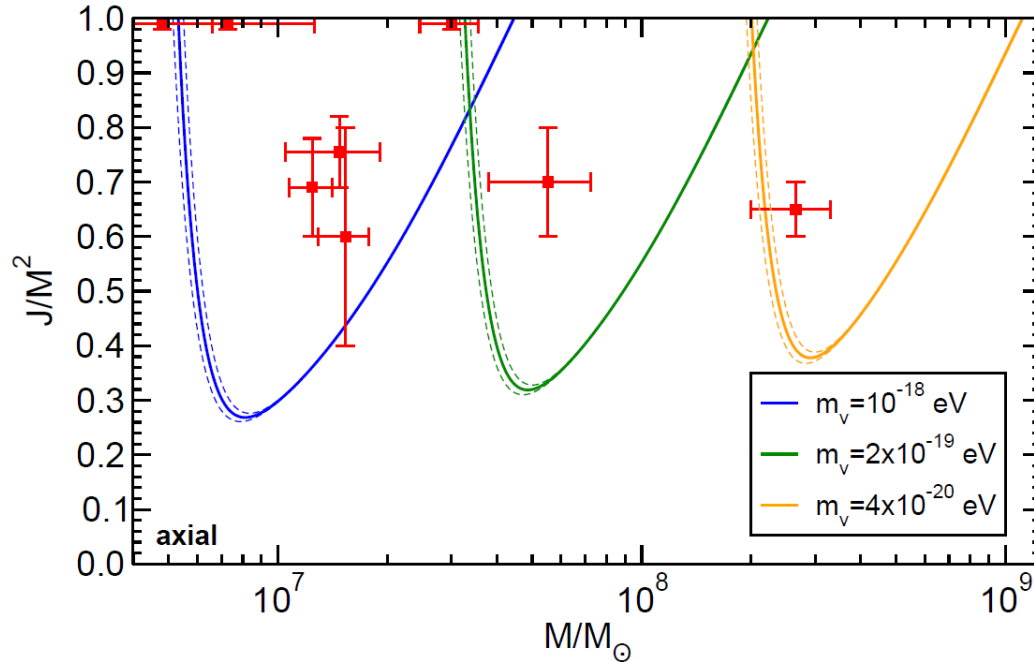


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Zel'dovich, Pis'ma Zh. Eksp. Teor. Fiz. 14 (1971), *Damour et al*, Lett.Nuovo Cim. 15 (1976) ;
Cardoso and Dias, Phys.Rev. D70 (2004) ; *Brito, Cardoso and Pani*, in preparation (2015)

Bounding the “photon” mass

Pani et al PRL109, 131102 (2012)



Depend very mildly on the fit coefficient and on the threshold

*

$\tau_{\text{Salpeter}} \rightarrow$ timescale for accretion at the Eddington limit

Bounding the graviton mass

Brito et al PRD88, 023514 (2013)

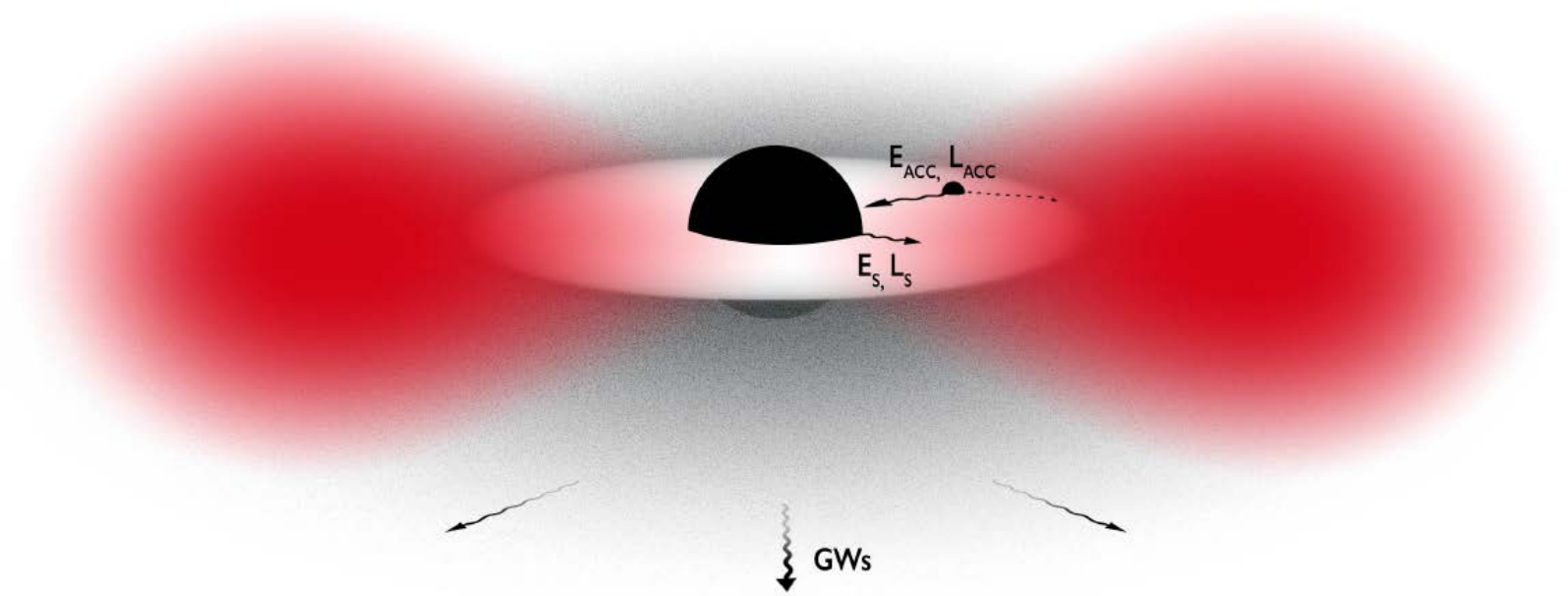
Bound on photon mass is model-dependent: details of accretion disks or intergalactic matter matter...but gravitons interact very weakly!

$$\square h_{\mu\nu} + 2\bar{R}_{\alpha\mu\beta\nu}h^{\alpha\beta} - \mu^2 h_{\mu\nu} = 0$$

<u>VALUE (eV)</u>	<u>DOCUMENT ID</u>	<u>COMMENT</u>
$\leq 6 \times 10^{-32}$	¹ CHOUDHURY 04	Weak gravitational lensing
● ● ● We do not use the following data for averages, fits, limits, etc. ● ● ●		
$< 5 \times 10^{-23}$	² BRITO 13	Spinning black holes bounds
$< 4 \times 10^{-25}$	³ BASKARAN 08	Graviton phase velocity fluctuations
$< 6 \times 10^{-32}$	⁴ GRUZINOV 05	Solar System observations
$> 6 \times 10^{-34}$	⁵ DVALI 03	Horizon scales
$< 8 \times 10^{-20}$	^{6,7} FINN 02	Binary pulsar orbital period decrease
	^{7,8} DAMOUR 91	Binary pulsar PSR 1913+16
$< 2 \times 10^{-29} h_0^{-1}$	GOLDHABER 74	Rich clusters
$< 7 \times 10^{-28}$	HARE 73	Galaxy
$< 8 \times 10^4$	HARE 73	2γ decay

Nonlinear effects are expected to be negligible!

Gravitational-wave emission and accretion?



Accretion:

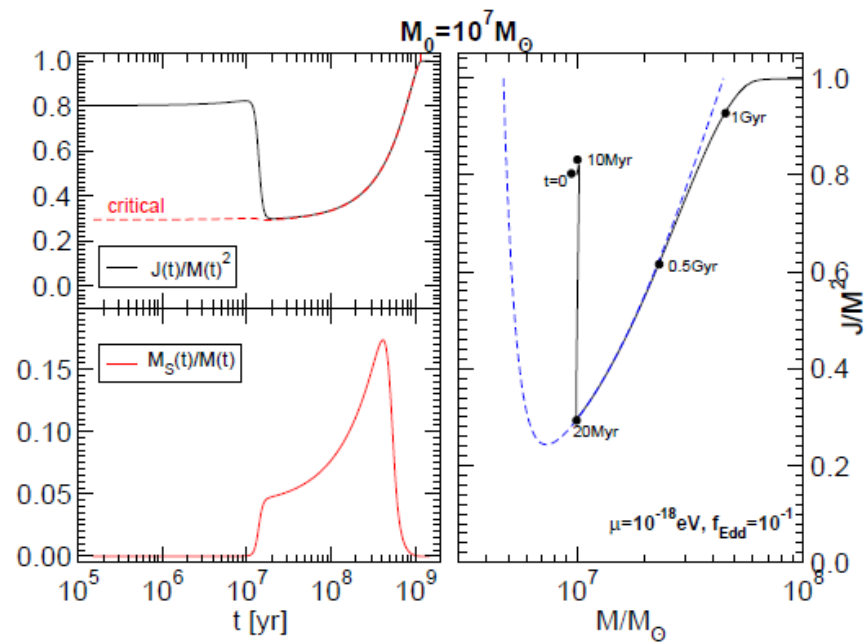
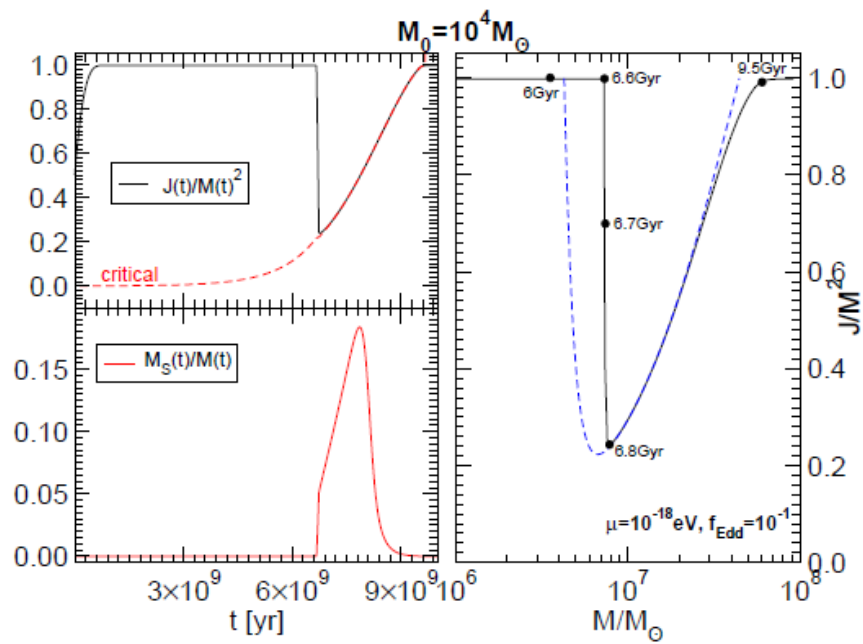
$$\dot{M}_{\text{ACC}} \equiv f_{\text{Edd}} \dot{M}_{\text{Edd}} \sim 0.02 f_{\text{Edd}} \frac{M(t)}{10^6 M_{\odot}} M_{\odot} \text{yr}^{-1}$$

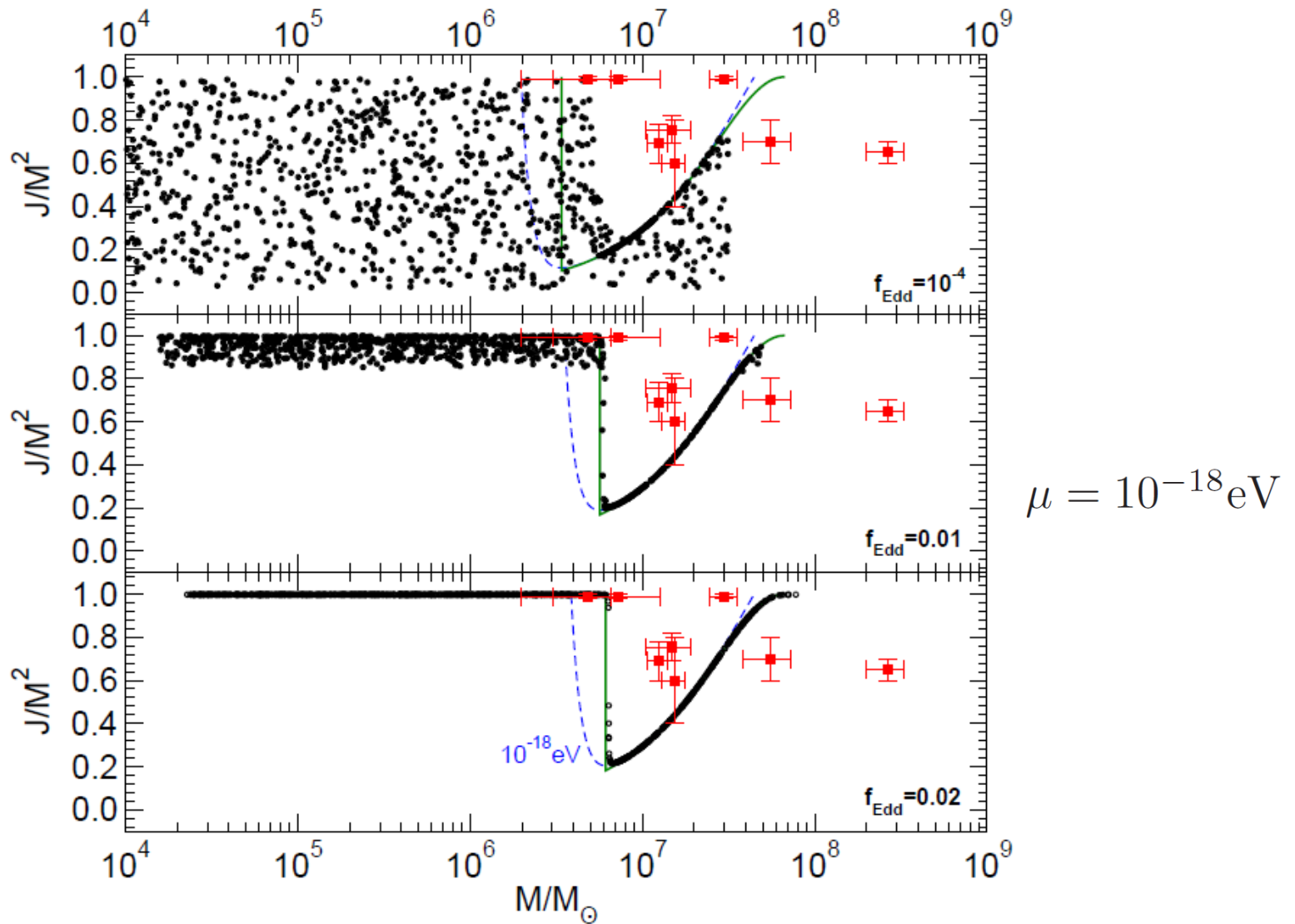
$$\dot{J}_{\text{ACC}} \equiv \frac{L(M, J)}{E(M, J)} \dot{M}_{\text{ACC}}$$

Gravitational-wave emission:

$$\dot{E}_{\text{GW}} = \frac{484 + 9\pi^2}{23040} \left(\frac{M_S^2}{M^2} \right) (M\mu)^{14}$$

$$\dot{J}_{\text{GW}} = \frac{1}{\omega_R} \dot{E}_{\text{GW}}$$





Random distributions 1000 BHs, with initial mass between $\log_{10} M_0 \in [4, 7.5]$ and $J_0/M_0^2 \in [0.001, 0.99]$ extracted at $t = t_F$, with t_F distributed on a Gaussian centered at $\bar{t}_F \sim 2 \times 10^9 \text{yr}$ with width $\sigma = 0.1\bar{t}_F$.

Strong field gravity is a fascinating topic

Fundamental fields, either in form of minimally coupled fields or under curvature couplings have a very rich and unexplored phenomenology: condensates outside BHs and compact stars act as gravitational-wave lighthouses, but can also act as dark matter.

Superradiant instabilities can provide strong constraints on masses of ultra-light bosons, turning black holes (and stars with dissipation channels) into effective particle detectors.

New BH solutions are possible stationary end-state of superradiant instability (see Herdeiro and Benone). Our results show that BHs grown out of Kerr are Kerr to a good precision

Thank you

